

# Cation Site and Anion Charge Ordering in Body-Centered Cubic-Packed Fullerides with Complex Counterions

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Two fullerides derived from  $\text{Na}_2\text{C}_{60}$  by ammonia coordination to the counterions are presented. Both structures adopt body-centered cubic-related anion packings. The sodium cations are extensively disordered in one phase, but their positional ordering in the second phase is correlated with the development of two distinct fulleride anion environments. Electrostatic considerations indicate that the two crystallographically distinct fulleride anions display charge separation.

## Introduction

The formation of complex counterions in fullerides by coordination of ammonia to intercalated metal cations has proved a versatile technique both for the expansion of fcc  $\text{A}_3\text{C}_{60}$  phases and the synthesis of novel fulleride structures.<sup>1</sup> The resulting expansion of the fulleride array can result in new electronic ground states, such as magnetic ordering in ammoniated  $(\text{NH}_3)\text{A}_3\text{C}_{60}$ <sup>2,3</sup> phases, which demonstrates the importance of electron correlation in the narrow band fullerides.  $\text{Na}_2\text{C}_{60}$  has previously been reported to afford rhombohedral  $(\text{ND}_3)_8\text{-Na}_2\text{C}_{60}$  and a body-centered cubic (bcc)-related  $(\text{NH}_3)_{5.3}\text{-Na}_2\text{C}_{60}$  phase.<sup>4</sup> In this paper we determine the structures of two bcc-packed ammoniated fullerides closely related to  $(\text{ND}_3)_{5.3}\text{Na}_2\text{C}_{60}$ , which differ in that occupancy of the interstitial cation sites is ordered in one case and disordered in the other. The cation site ordering, which is driven by the exceptional level of positional disorder found in the disordered phase, is accompanied by the development of two distinct fulleride anion environments. This is consistent with the formation of a charge-ordered array in which the crystallographically distinct anions carry different charges. In view of the importance of such valency ordering in controlling the electronic properties of oxides, and the increasing indications that fullerides are strongly correlated electronic systems, this is an important observation and further reinforces the important role of ammonia coordination to metal cations in discovering new fulleride ground states.

As is common in fulleride science, both structures were determined from powder diffraction data from heavily disordered materials. Both structures are based on body-centered cubic (bcc) anion arrays in which there are six pseudo-tetrahedral interstitial sites per fulleride which are occupied by ammonia molecules. This results

in many possible locations for two metal cations (the most favorable sites, occupied in  $(\text{NH}_3)_6\text{Na}_3\text{C}_{60}$ ,<sup>5,6</sup> would accommodate twelve sodium cations per formula unit if fully occupied) and extensive positional disorder. In this paper we show that ordering of the cations over these interstitial sites is possible, but solution of the resulting structure using conventional diffraction methods proves challenging. One method of overcoming this problem is the application of maximum entropy methods (MEM).<sup>7–10</sup> In this approach, rather than performing a Fourier transformation of the  $F_{\text{obs}}$  structure factors as obtained from the Rietveld refinement by apportioning the intensity of overlapping observed reflections according to the calculated structure factors, a density distribution is estimated and the information entropy contained in the structure factors is maximized within the errors of the observed data. Thus, MEM can be seen in some ways as the opposite of a Fourier synthesis, as an initial density is fitted to the observed structure factors (it is customary to start from a homogeneous density distribution). As unobserved high  $Q$  ( $Q = 2\pi/d$ ) peaks therefore do not effect the density, accurate electron densities can be obtained from limited numbers of integrated intensities.<sup>10</sup> The use of MEM alone still maintains the  $F_{\text{obs}}$  as obtained from the initial extraction and is therefore biased toward the model used to extract the data. The results of an MEM calculation can, however, be used to provide  $F_{\text{calc}}$  for a second extraction, which will be less biased by the original model. Repeated extractions and MEM calculations should therefore yield the true scattering density. This cycling of the MEM results has been implemented by Izumi in RIETAN-2000 who has applied the REMEDY cycles to a number of diffraction problems.<sup>11,12</sup>

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**Table 1. Refined Positional Parameters of  $(ND_3)_6Na_2C_{60}$  in Space Group  $Im\bar{3}$  (number 204) ( $a = 12.151(3)$  Å)**

name	<i>x</i>	<i>y</i>	<i>z</i>	100 U/Å <sup>2</sup>	fractional occupancy	site
C1	0.05667(40)	0.00000	0.28850(32)	4.3 (2)	1	24g
C2	0.11788(21)	0.09844(27)	0.25287(24)	4.9 (1)	1	48h
C3	0.05696(19)	0.18544(23)	0.21382(22)	5.6(1) 1	1	48h
Na4	0.31978(33)	0.34151(29)	0.00000	10(4)	0.166	24g
N5	0.50000	0.2837(7)	0.00000	8.2(2)	1	12e
D7	0.5589(13)	0.2241(14)	0.0341(7)	13.9(6)	0.25	48h
D8	0.4500(12)	0.3283(13)	0.0608(9)	13.9(6)	0.25	48h
D9	0.4560(15)	0.2567(14)	-0.0743(12)	13.9(6)	0.25	48h

## Experimental Section

$(ND_3)_6Na_2C_{60}$  was prepared from  $(ND_3)_8Na_2C_{60}$  (synthesized as described in ref 4<sup>4</sup>) by equilibration of the initial  $(ND_3)_8Na_2C_{60}$  phase with 0.05 atm of  $ND_3$  vapor at 100 °C.  $(NH_3)_{5.3}Na_2C_{60}$  was prepared as described previously<sup>4</sup> by equilibration of  $Na_2C_{60}$  with ammonia gas. The composition was determined volumetrically. Neutron powder diffraction data on  $(ND_3)_6Na_2C_{60}$  were collected at ambient temperature on the POLARIS powder diffractometer<sup>13</sup> at the ISIS spallation neutron source. The relevant neutron scattering lengths are  $C = 0.665$ ,  $D = 0.667$ ,  $N = 0.930$ , and  $Na = 0.363 \times 10^{-12}$  cm. Synchrotron X-ray powder diffraction data were recorded at ambient temperature at station 9.1 of the Daresbury Synchrotron Radiation Source. Data analysis was performed using the GSAS refinement and REMEDY (a combination of RIETAN-2000<sup>12</sup> and MEED<sup>9</sup>) maximum entropy analysis packages. Room temperature <sup>13</sup>C MAS NMR data were recorded on  $(NH_3)_{5.3}Na_2C_{60}$  on a Bruker MSL200-S NMR spectrometer with an Oxford Instruments 4.7 T wide bore (98 mm) superconducting solenoid magnet equipped with an Aspect 2000 data system. Spectra were recorded at 50.32 MHz using a Bruker Z32 HB-WB-73A MAS7DB probe. All samples for MAS NMR were packed into 5-mm Kel-F capped inserts which were placed into 7-mm zirconia rotors with Kel-F caps – this produced an airtight seal. The spinning speed was 4.2 kHz, and the recycle delay was 20 s.

## Results

The ammonia sorption isotherms of  $Na_2C_{60}$  display considerable hysteresis, and this is reflected in the subtle structural differences between the two phases discussed here. Both phases give X-ray powder patterns which can be predominantly indexed with the bcc unit cell of  $(NH_3)_6Na_3C_{60}$ , but differ in that  $(NH_3)_{5.3}Na_2C_{60}$  (prepared directly from  $Na_2C_{60}$ ) displays superstructure reflections characteristic of a primitive trigonal superstructure derived from a trigonal distortion of the bcc subcell adopted by  $(ND_3)_6Na_2C_{60}$ , prepared by deamination of  $(ND_3)_8Na_2C_{60}$ . We discuss the structure of  $(ND_3)_6Na_2C_{60}$  derived by neutron powder diffraction initially because it is the disordered precursor of the ordered  $(NH_3)_{5.3}Na_2C_{60}$ . Le Bail extraction of the 3 banks of neutron data converged smoothly to  $\chi^2 = 2.53$ . Analysis of the room-temperature diffraction pattern began with the  $(NH_3)_6Na_3C_{60}$  model.<sup>5,6</sup> Refinement of the N and Na fractional occupancies indicated that the N site was fully occupied, corresponding to a composition  $(ND_3)_6Na_2C_{60}$ , and thus the site occupancies were fixed at the values determined by this composition in the subsequent analysis. Difference Fourier analysis to

**Table 2. Bond Lengths (Å) and Angles (°) in  $(ND_3)_6Na_2C_{60}$  Refined from Neutron Powder Diffraction**

C1–C1 (×6)	1.38(1)
C2–C3 (×24)	1.375(3)
C1–C2 (×24)	1.474(5)
C2–C3 (×24)	1.511(4)
C3–C3 (×12)	1.384(5)
C1–C1–C2 (×2)	120.3(2)
C2–C1–C2	108.5(4)
C1–C2–C3	117.0(3)
C1–C2–C3	106.1(3)
C3–C2–C3	118.9(3)
C2–C3–C2	120.2(3)
C2–C3–C3	122.6(2)
C2–C3–C3	109.5(2)
Na4–N5	2.299(2)
Na4–N5	2.300(5)
N5–Na4–N5	140.9(3)
N5–D7	1.099(5)
D7–N5–D8	115.5(6)

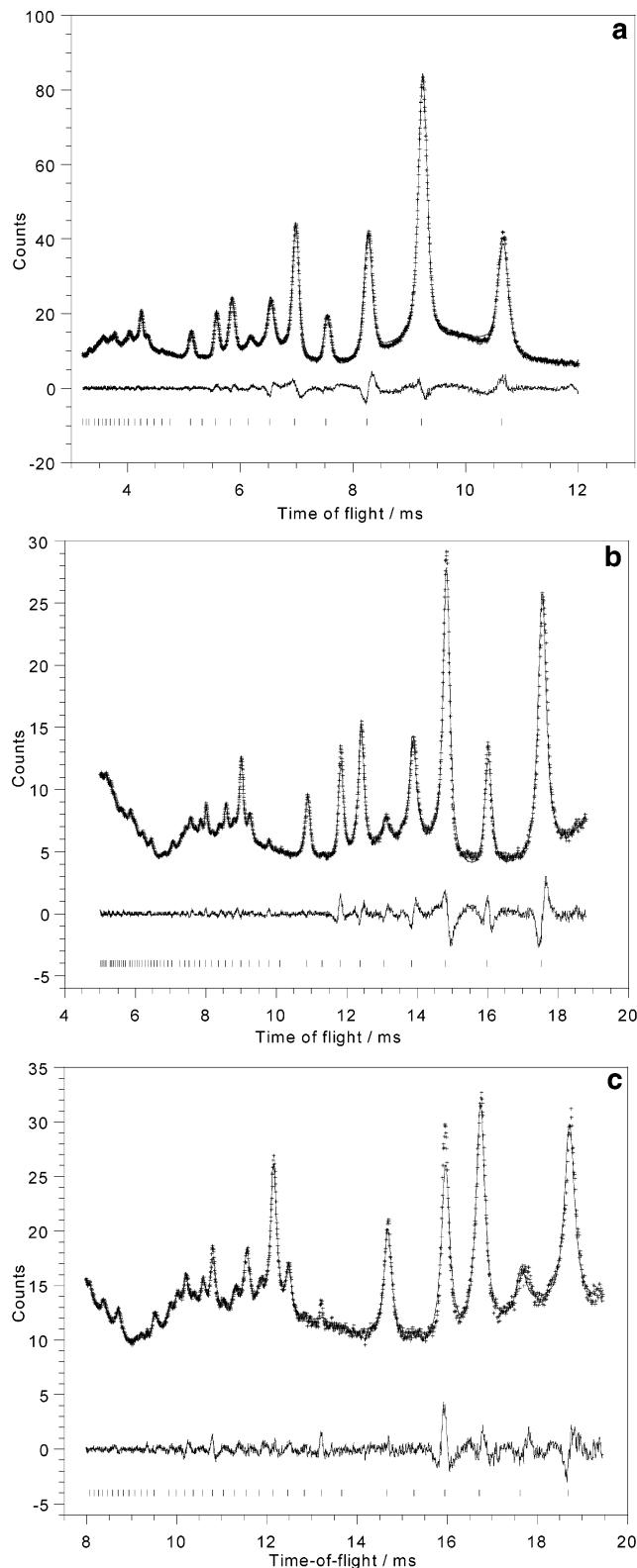
locate the D positions did not yield discrete well-defined maxima, ascribable to the orientational disorder evidenced in the diffraction pattern in which Bragg peaks are not discernible below 1.2 Å. Data collected at 5 K indicate that this positional disorder becomes frozen-in on cooling. A rigid body (initial geometry N–D = 1.12 Å, D–N–D = 109°) model was therefore used to describe the deuterium positions in this phase. Location of the rigid body  $ND_3$  molecules on the 12e ( $1/2$ ,  $y$ , 0) special positions and difference Fourier analysis lead to the location of Na on 24g ( $x$ ,  $y$ , 0) as in  $(NH_3)_6Na_3C_{60}$ . Both refinement and analysis of difference Fourier maps indicated that there was no N within the vicinity of this site. Refinement of the rigid body rotation angles, N–D bond length and D–N–D angle, site occupancy factors, displacement parameters, and positional parameters for non- $ND_3$  atoms lead to the model presented in Table 1 and the fit shown in Figure 1. The structure is shown in Figure 2a (The ORTEP representation is provided in Figure S1 of the Supporting Information). The  $\chi^2$  value of 2.7 is acceptably close to the Le Bail extraction, indicating that the remaining inadequacies are due to difficulties in describing the peak shape. Well-determined bond-distances and angles are given in Table 2. Refinement in space group  $Im\bar{3}m$ , in which the  $C_{60}$  anion positions are disordered over two orientations related by 90° rotation about [001], was unstable and did not converge below  $\chi^2 = 5.0$ .

The  $C_{60}^{2-}$  anions adopt  $m\bar{3}$  symmetry standard orientations in the bcc array. The  $ND_3$  molecules are orientationally disordered within distorted tetrahedral sites in this  $C_{60}^{2-}$  array (Figure 2a), with closest C–D contacts of 2.674 Å (C1–D3) and 2.57 Å (C2–D2) (as represented in Figure S2 of the Supporting Information) and several other contacts closer than 2.95 Å seen in solid benzene. This is indicative of N–D...π interactions playing a role in determining the orientation of the  $ND_3$

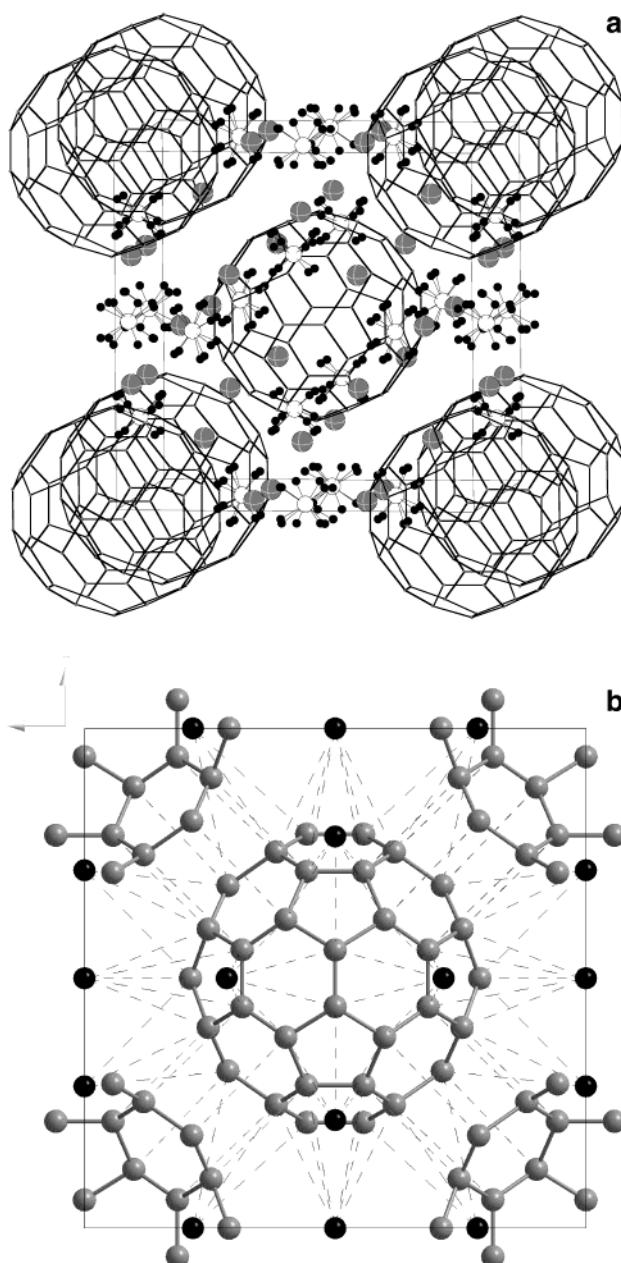
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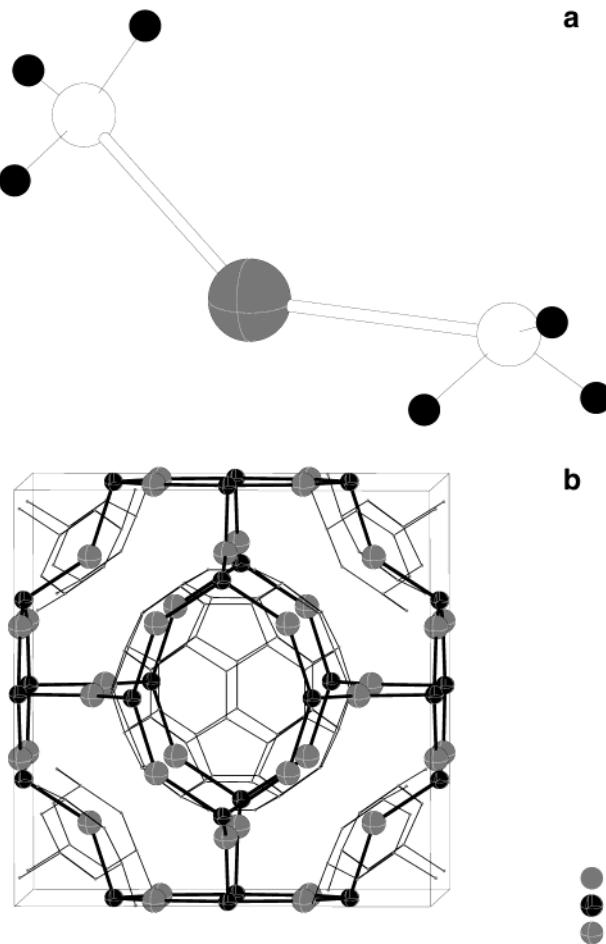


**Figure 1.** Rietveld fit to time-of-flight neutron powder diffraction data from  $(ND_3)_6Na_2C_{60}$ . (a)  $35^\circ$  detector bank ( $1.45 \leq d/\text{\AA} \leq 5.6$ ) (b)  $90^\circ$  detector bank ( $1.1 \leq d/\text{\AA} \leq 4$ ) (c)  $145^\circ$  detector bank ( $1.29 \leq d/\text{\AA} \leq 3.2$ ). The observed data are shown as points, the calculated fit is the solid line, and the difference is plotted below. The tick marks indicate the location of the Bragg peaks.  $\chi^2 = 2.71$  (86 variables, 7071 data points) for the fit to all three data banks simultaneously. The agreement indices for the individual banks are as follows:  $35^\circ R_w = 4.09\%$ ,  $R(F^2) = 1.88\%$  for 107 reflections;  $90^\circ R_w = 2.47\%$ ,  $R(F^2) = 2.47\%$  for 263 reflections;  $145^\circ R_w = 2.52\%$   $R(F^2) = 0.96\%$  for 154 reflections. The refined parameters are given in Table 1.



**Figure 2.** (a) Refined structure of  $(ND_3)_6Na_2C_{60}$  (N unshaded, D black, Na gray crossed spheres). The Na sites are  $1/6$  occupied and thus the structure is considerably positionally disordered. (b) The environment of a  $C_{60}^{2-}$  anion in  $(ND_3)_6Na_2C_{60}$  defined by contacts to the  $ND_3$  molecules and eight neighboring  $C_{60}^{2-}$  anions along the  $\langle 111 \rangle$  directions, which cap the eight hexagonal faces not involved in  $ND_3-C_{60}$  interactions.

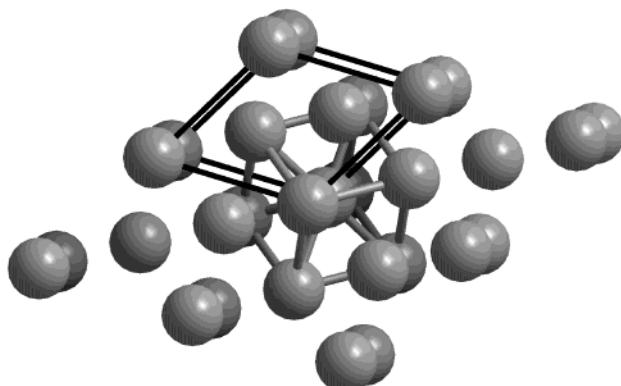
units. Each  $C_{60}$  is surrounded by 24  $ND_3$  units: of the 32 faces, all 12 pentagonal, and 12 of the 20 hexagonal faces are capped with  $ND_3$  molecules. The eight fulleride neighbors of a given  $C_{60}^{2-}$  anion are located 10.52 Å away along the body diagonals of a cube, and the remaining eight hexagonal faces of the anion are aligned with these directions, remaining uncapped by ammonia molecules (Figure 2b). The close C–D contacts arise from  $ND_3$  neighbors located over hexagonal faces, and allow the  $ND_3$  molecules to link the  $C_{60}^{2-}$  anions via these interactions, as the close C...D2 and C...D3 contacts for a given  $ND_3$  are to neighboring anions. None of the D atoms are located over hexagon or pentagon centroids as in  $(ND_3)_8Na_2C_{60}$ ,<sup>4</sup> and this may account for



**Figure 3.** (a) Coordination environment of the  $\text{Na}^+$  cations in  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$ . The partial occupancy of the cation sites and orientational disorder of the anion means that other local cation environments are possible. The D atoms are in black, N are unshaded, and Na are gray crosses. (b) The extended network of contacts between the fully occupied  $\text{ND}_3$  (dark crossed spheres) sites and the  $\frac{1}{6}$  occupied Na sites (gray crossed spheres). As each  $\text{ND}_3$  makes only four contacts to Na sites (Figure S3 in the Supporting Information), there are  $\text{ND}_3$  molecules which do not coordinate to  $\text{Na}^+$ .

the absence of orientational ordering of the ammonia molecules in the present case.

The  $(\text{ND}_3)_2\text{Na}^+$  complexes (Figure 3) refine with an N–Na–N angle of  $140.9(3)^\circ$ ; the Na–C3 contact of  $2.356(5)$  Å is shorter than the radius sum of  $2.7$  Å and there may be local relaxation of the anion orientation in the vicinity of the 14% occupied Na site. As a result, C3 has the largest displacement parameter of any of the C atoms. The  $\text{ND}_3$  site is located equidistant from four partially (3) occupied Na sites, and the D positions thus represent an average over four possible orientations of the  $\text{ND}_3$  unit, together with  $\text{ND}_3$  molecules uncoordinated to Na cations (such  $\text{NH}_3$  species are observed in the structures of  $\text{Ba}(\text{NH}_3)_7\text{C}_{60} \cdot \text{NH}_3^{14}$  and similar phases). This prevents the definitive identification of  $\text{ND}_3$  orientations associated with a specific  $\text{Na}(\text{ND}_3)_2$  complex, with Figure 3a representing the most chemically sensible local description. More complex models, with multiple  $\text{ND}_3$  orientations, were tested but could not be converged satisfactorily. A model involving



**Figure 4.** Relationship between the sodium site-disordered  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  (cube delineated by gray solid lines) and sodium site-ordered  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  unit cell (black lines), viewed close to the [001] direction of the trigonal  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  cell. Fulleride anions in the pseudo-bcc array are represented as spheres and the eight nearest neighbors in the bcc cell are connected by gray solid lines.

disorder of the  $\text{ND}_3$  units over the vertexes of a cube converged to  $\chi^2 = 3.8$ , indicating that the observed orientational disorder of the  $\text{ND}_3$  units does not correspond to rotational diffusion. The network of  $\text{Na} \dots \text{N}$  contacts is shown in Figure 3b.

$(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  adopts a trigonal structure which is derived from a rhombohedral distortion ( $c/a = 0.609$ ) of  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  ( $c_{\text{trig}}/a_{\text{trig}} = 0.6124$  for the undistorted bcc cell) (Figure 4). This corresponds to compressing the bcc unit cell along the  $\langle 111 \rangle$  direction. The diffraction pattern shows no systematic absences, and so given the clear structural relationship to  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  further analysis was conducted in  $P\bar{3}$  which is the only centric trigonal group which can be derived from maximal nonisomorphic subgroups of  $I\bar{m}\bar{3}$ . Initial models were constructed by refining the data in  $I\bar{m}\bar{3}$  using the Na and N positions found in the disordered  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  structure above, followed by reduction of symmetry to  $P\bar{3}$ . Neutral atom scattering factors were used for Na and C,  $\text{NH}_3$  was approximated by 1.42 times the scattering factor for nitrogen to allow for the electron density due to the protons. This corresponds to assigning the symmetry lowering from cubic to trigonal to the ordered occupancy of the disordered metal and ammonia sites in the cubic space group. Refinement of this model yields some order of the sodium and nitrogen sites. However, the resultant model was not physically reasonable because the Na–Na and N–N nonbonded contacts are too short given the resulting occupancies. We therefore concluded that the data did not allow a unique and physically acceptable refinement of this highly disordered system by this approach, and the absence of diffraction data at high angle defeated efforts to use conventional Fourier methods to resolve the disorder.

Therefore, this model was used as a starting point for REMEDY<sup>15,16</sup> cycles using a combination of RIETAN-2000 and the MEED/MEVIUS programs developed by Sakata and co-workers.<sup>9</sup> Initially  $F_{\text{obs}}$  were

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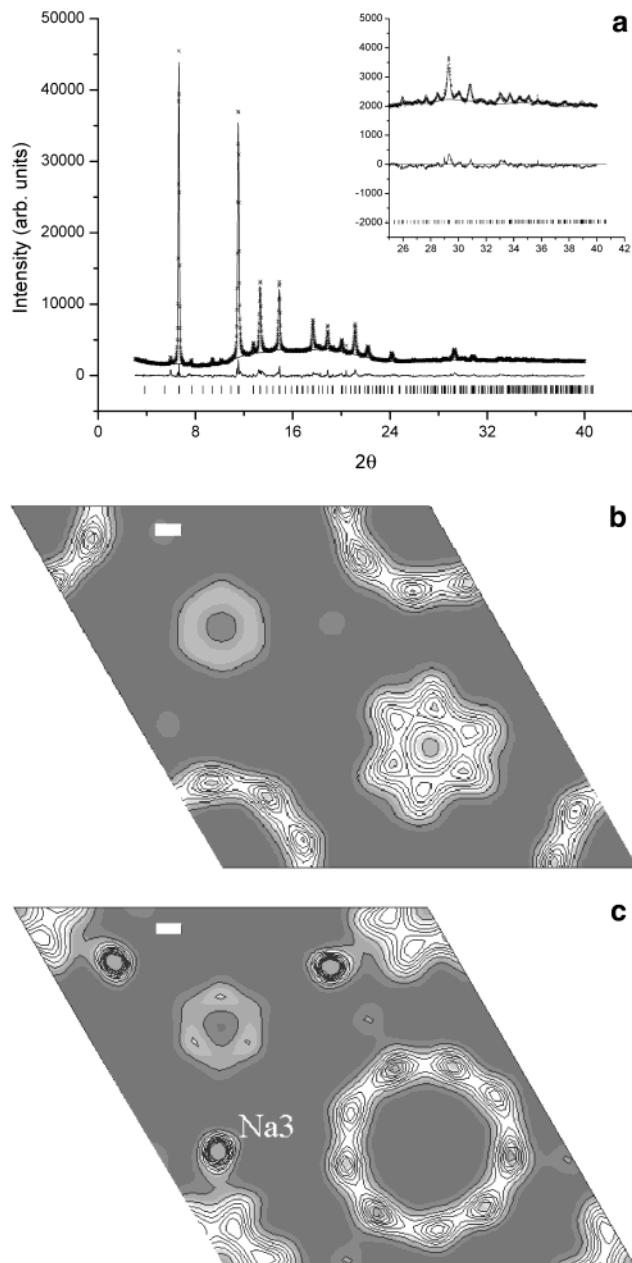
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extracted on the basis of the model described in the previous paragraph. These values were then used to calculate an electron density map, using the method first employed by Collins<sup>7</sup> as implemented by Sakata and co-workers.<sup>9,10</sup> The electron density in the unit cell is divided into a number of voxels (in this case  $64 \times 64 \times 64$ ) and the electron density refined against  $F_{\text{obs}}$  starting from an initial featureless map. The resulting density map is used to calculate a set of  $F_{\text{calc}}$  values which are used to extract a further set of  $F_{\text{obs}}$ . The process is repeated in order to remove the bias of the initial model. No further improvement in the fit was obtained after 4 cycles ( $R(F^2)$  improved from 9.58% to 2.05%, as shown in Figure 5a). The electron density map thus generated clearly showed the different setting angles of the symmetry-independent sets of fulleride anions (Figure 5b) and allowed the location of five sites not associated with the  $\text{C}_{60}$  anions, all of which were general positions (6g). Two of these sites were in close proximity ( $< 2 \text{ \AA}$ ), and were thus refined as being partially occupied (with the sum of the occupancies fixed at 1.42, which approximates to 1  $\text{NH}_3$ , *vide supra*). Given the identical electron counts of the  $\text{Na}^+$  and  $\text{NH}_3$  units, assignment of the site occupancy to chemical species was made using chemical bonding and coordination number considerations. The sodium was chosen to be on the site with the most reasonable distances to carbon (Figure 5c), which is also coordinated by two of the other non- $\text{C}_{60}$  sites. The remaining sites were assumed to be fully occupied by  $\text{NH}_3$ , leading to a final refined composition of  $(\text{NH}_3)_{5.73(4)}\text{Na}_2\text{C}_{60}$ . Given the disorder present in the material, this is acceptably close to the composition determined from the ammonia sorption isotherm.

An important feature of the symmetry lowering from cubic to trigonal is that two crystallographically inequivalent  $\text{C}_{60}^{2-}$  sites are present: one is located at the *1a* position (0,0,0; site symmetry  $\bar{3}$ ) and the second is located on  $2d(1/3, 2/3, z, z \sim 2/3, \text{site symmetry } 3)$ . Because of the absence of diffraction at high angle, both of these anions were modeled as rigid bodies with only the rotation angle of each anion around the 3-fold axis and the *z* displacement of the *2d* anion refined. Initial orientations for the  $\text{C}_{60}^{2-}$  anions were obtained from the MEED map, and the setting angles were then refined to values close to these initial ones. Final Rietveld fits, which employed an overall thermal parameter, are shown in Figure 6a, and the resulting  $\text{C}_{60}$  orientations are shown in Figure 6b. The refined parameters are given in Table 3, and the resulting bond lengths and angles are provided in Table 4.

$(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  gives a single peak in the room temperature  $^{13}\text{C}$  MAS spectrum at 184 ppm, compared to the shift of 172 ppm found for  $\text{Na}_2\text{C}_{60}$ <sup>17</sup> (Figure 7). The full width at half-height is 115 Hz, compared with the value of 28 Hz found for the  $\text{Na}_2\text{C}_{60}$  starting material. This considerable enhancement in the breadth of the  $^{13}\text{C}$  MAS NMR resonance is consistent with the development of two similar but distinct  $\text{C}_{60}$  sites as required by the space group symmetry.

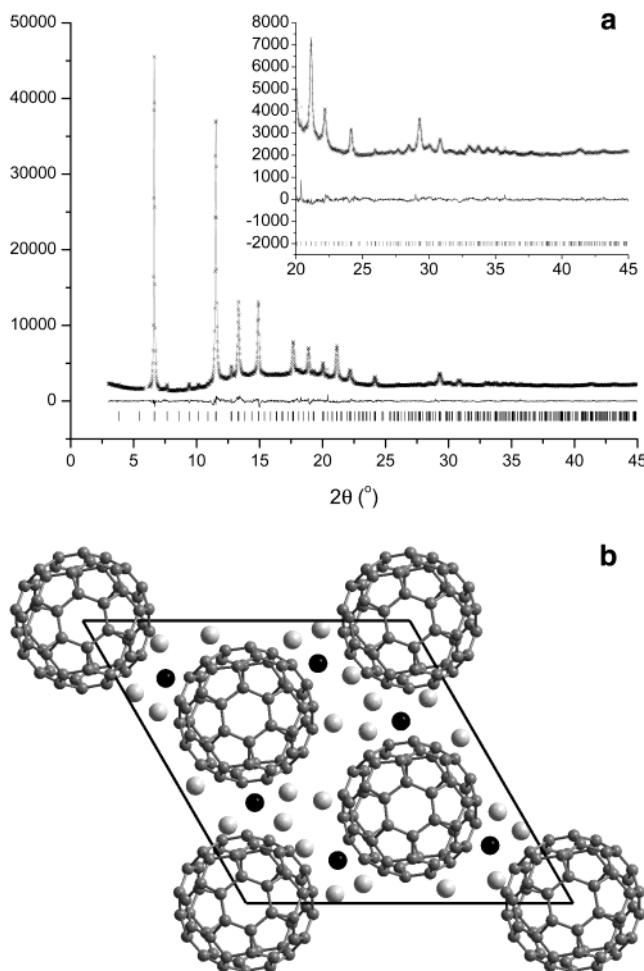
The sodium is unsymmetrically bound by the fully occupied  $\text{NH}_3\text{b}$  site at 2.23(3)  $\text{\AA}$  and the approximately



**Figure 5.** (a) REMEDY fit to the diffraction pattern of  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$ . (b) Sections of the *ab* plane at *z* = 0.05 of the MEED/REMEDY electron density of  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  showing the orientation of the  $\text{C}_{60}$  on the *2d* position (contours are marked at 0.3  $\text{e\AA}^{-3}$ , scale bar is 1  $\text{\AA}$ ). (c) Sections of the *ab* plane at *z* = 0.28 of the MEED/REMEDY electron density of  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  showing the location of the  $\text{Na}3$  site (contours are marked at 0.3  $\text{e\AA}^{-3}$ , scale bar is 1  $\text{\AA}$ ).

half occupied  $\text{NH}_3\text{c}$  and  $\text{NH}_3\text{d}$  sites at 3.04(7) and 2.82(8)  $\text{\AA}$ , respectively. The sodium site is approximately equidistant between two 6:6 ring junctions of the two different  $\text{C}_{60}$  sites (Figure 8). Figure 8b,c shows a comparison of location of the sodium site relative to the six sites generated by symmetry lowering from the positionally disordered  $\text{Im}\bar{3}$  phase. The distances between the occupied site in the ordered phase and the parent disordered sites vary between 4.334 and 1.485  $\text{\AA}$ , with the two closest sites at 1.485 and 1.701  $\text{\AA}$ . The sodium cation has moved from the centroid of these sites to a position between the two  $\text{C}_{60}^{2-}$  anions which coordinate to it.

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**Figure 6.** (a) Rietveld fit to synchrotron X-ray powder diffraction data from  $(ND_3)_{5.3}Na_2C_{60}$ . The observed data are shown as points, the calculated fit is the solid line, and the difference is plotted below. The tick marks indicate the location of the Bragg peaks ( $\chi^2 = 3.65$ ,  $R_{wp} = 2.83\%$ ,  $R_p = 2.29\%$ ,  $R(F^2) = 6.48\%$ .) Space group  $P\bar{3}$  (no. 147)  $a = 17.227(1)$  Å,  $c = 10.500(1)$  Å,  $V = 2698.9(5)$  Å $^3$ . The refined parameters are given in Table 3 and the lengths are provided in Table 4. (b) Refined setting angles of the two independent  $C_{60}^{2-}$  anions, viewed along [001]. Na cations are represented as black and  $NH_3$  molecules are shown as light spheres.

The ordering of the sodium positions is coupled crystallographically to the development of the two distinct fulleride environments. The occupied sodium sites form an octahedron around the  $C_{60}$  on site  $1a$  (the Na sites are 6.101 Å from the anion centroid). The other  $C_{60}$  site ( $2d$ ), which relaxes by 0.355(3) Å along  $c$  toward the  $1a$  site, also has six Na nearest neighbors. However, the sodium sites are now arranged in a distorted trigonal prism. There are three short Na-centroid distances (6.381 Å) forming one triangular face of the trigonal prism, and three long distances (7.705 Å). The arrangements of the sodiums and ammonias around the two different  $C_{60}$  sites are shown in Figure 9. The mean Na– $C_{60}$  distance is considerably longer to the  $2d$   $C_{60}$ , consistent with different anion charges. (Assuming that the interactions between the  $Na^+$  and  $C_{60}$  anions are purely electrostatic, the distances between the  $C_{60}$  centroids and the cations will depend on the relative charges on the  $C_{60}$  anions, and based on this simple estimate a 15% charge difference is indicated.) This is the first example of the development of two different

**Table 3. Refined Parameters for  $Na_2(NH_3)_{5.3}C_{60}^a$**

	<i>x</i>	<i>y</i>	<i>z</i>	fractional occupancy
C1a	0.17396(28)	0.1038(21)	0.2324300	1
C2a	0.2217(12)	0.0408(34)	0.0548110	1
C3a	0.09501(24)	0.0618(11)	0.3120690	1
C4a	0.1913(14)	0.1789(17)	0.1549790	1
C5a	0.1906(10)	0.0347(29)	0.1826350	1
C6a	0.0342(13)	0.09512(23)	0.3133760	1
C7a	0.1265(25)	0.2097(4)	0.1529670	1
C8a	0.0508(24)	0.1693(6)	0.2294000	1
C9a	0.1223(30)	0.23948(4)	0.0242370	1
C10a	0.2229(12)	0.1848(22)	0.0248410	1
C1b	0.54143(34)	0.7399(19)	0.5398(8)	1
C2b	0.12527(34)	0.5934(19)	0.8457(8)	1
C3b	0.2743(16)	0.49561(30)	0.4634(8)	1
C4b	0.16232(30)	0.5546(13)	0.9221(8)	1
C5b	0.2762(8)	0.7055(9)	0.3794(8)	1
C6b	0.3904(7)	0.6279(9)	0.1.0062(8)	1
C7b	0.2269(21)	0.7991(19)	0.6685(8)	1
C8b	0.4398(21)	0.5342(19)	0.7170(8)	1
C9b	0.3087(19)	0.8288(12)	0.8754(8)	1
C10b	0.1465(8)	0.6421(19)	0.5102(8)	1
C11b	0.1607(12)	0.6898(21)	0.8478(8)	1
C12b	0.3565(21)	0.8625(8)	0.5377(8)	1
C13b	0.3912(7)	0.6288(9)	0.3806(8)	1
C14b	0.2755(7)	0.7045(9)	1.0049(8)	1
C15b	0.1563(15)	0.7166(23)	0.7176(8)	1
C16b	0.5104(15)	0.6167(23)	0.6679(8)	1
C17b	0.5508(9)	0.6957(23)	0.7475(8)	1
C18b	0.1159(9)	0.6376(23)	0.6380(8)	1
C19b	0.2368(14)	0.7449(15)	0.9252(8)	1
C20b	0.4116(15)	0.84131(10)	0.4603(8)	1
Na3	0.2044(11)	0.3560(15)	0.2824(19)	1
NH <sub>3a</sub>	0.0519(21)	0.4151(21)	0.2101(31)	1.35(2)
NH <sub>3b</sub>	0.3238(18)	0.3923(19)	0.1597(31)	1.29(3)
NH <sub>3c</sub>	0.194(4)	-0.081(5)	0.510(6)	0.64(4)
NH <sub>3d</sub>	0.287(4)	0.031(5)	0.511(7)	0.78(4)

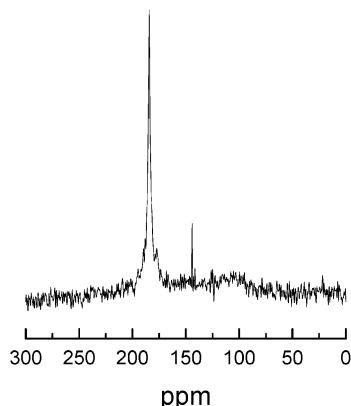
<sup>a</sup> All the atoms occupy the  $6g$  positions. The  $NH_3$  groups were refined using the N form factor, and thus occupancies should be divided by 1.42 to yield the fractional occupancy of  $NH_3$ , as discussed in the text. The  $C_{60}$  anions were refined as rigid bodies using two parameters for the  $C_{60}$  on the  $2d$  site (translation along  $c$  and rotation around the 3-fold axis) and one parameter for the  $C_{60}$  on the  $1a$  site (rotation around the 3-fold axis).

**Table 4. Refined Bond Lengths (Å) and Angles (°) in  $(NH_3)_{5.3}Na_2C_{60}$**

Na3–C7a	2.57(2)
Na3–C8a	3.02(3)
Na3–C3b	2.82(2)
Na3–C1b	3.25(3)
Na3–NH <sub>3b</sub>	2.23(2)
Na3–NH <sub>3c</sub>	3.04(7)
Na3–NH <sub>3d</sub>	2.82(8)
C8a–Na3–C3b	169.3(33)
C8a–Na3–NH <sub>3d</sub>	85.8(26)
C3b–Na3–NH <sub>3d</sub>	102.4(4)

anion sites in a fulleride in which the anions are in electronic contact, and may indicate an uneven distribution of the  $t_{1u}$  electrons between the distinct anions, producing a real-space charge-ordered state. This is also reflected in the  $C_{60}$ – $C_{60}$  closest contacts. The  $C_{60}$  ( $1a$ ) is surrounded by a regular bicapped octahedron of  $C_{60}$ 's ( $2 \times 10.497$  Å to  $1a$  and  $6 \times 10.443$  Å to  $2d$ ), whereas the  $C_{60}$ – $C_{60}$  closest contacts around the  $2d$  site are significantly distorted ( $3 \times 10.443$  Å to  $1a$  and  $2 \times 10.497$  Å,  $3 \times 10.782$  Å to  $2d$ ).

The three nitrogen sites shift significantly less than the sodium sites when compared with the disordered parent material, and the reduced displacement of these neutral species is further evidence of coupling of the cation positional and anion charge ordering. The non-

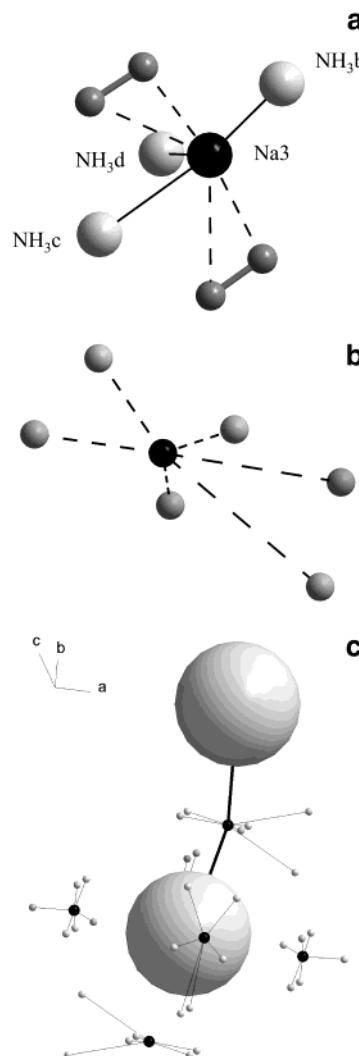


**Figure 7.**  $^{13}\text{C}$  MAS NMR of  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$ . (The sample on which NMR data were recorded was a larger preparation than that used for diffraction data collection and thus contained a small amount of unreacted  $\text{C}_{60}$  evident as a sharp line at 143 ppm.)

coordinated  $\text{NH}_3\text{a}$  moves 0.824 Å toward the center of the pseudo-tetrahedral cavity at  $\frac{1}{2}$ , 0,  $\frac{1}{4}$  and the coordinated  $\text{NH}_3\text{b}$  moves 0.557 Å toward the occupied Na site. The remaining site splits and both components move toward the two adjacent sodium sites (by 1.521 and 0.418 Å). The presence of ammonia not coordinated to sodium in this site-ordered phase indicates the favorable nature of this site in a bcc fulleride array, and gives support to the inference of such noncoordinating ammonia in the site-disordered  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  phase.

## Discussion

Both phases discussed here are derived from a body-centered packing of the fulleride anions. The reactivity with ammonia of fcc  $\text{Na}_2\text{C}_{60}$ , in which the sodium cations are ordered on the tetrahedral sites, is driven by the need to coordinate all of the small sodium cations to ammonia, which favors more expanded bcc-based structures to allow the  $\text{Na}(\text{NH}_3)_x$  complexes to form. In contrast, the fcc phases  $\text{Na}_2\text{CsC}_{60}$  and  $\text{K}_3\text{C}_{60}$  afford  $(\text{NH}_3)_4\text{Na}_2\text{CsC}_{60}$ <sup>18</sup> and  $(\text{NH}_3)_4\text{K}_3\text{C}_{60}$ <sup>19</sup> with  $(\text{NH}_3)_4\text{Na}^+$  and  $[\text{K}-\text{NH}_3]^+$  units on the octahedral sites, while the tetrahedral sites are occupied by uncomplexed cations. The bcc array contains large quasi-tetrahedral interstitial sites which are occupied by the cations in  $\text{K}_6\text{C}_{60}$ .<sup>20</sup> Although neutral fcc-packed  $\text{C}_{60}$  does not take up ammonia, these large interstitial sites in the bcc array are favorable locations for the ammonia molecules and are occupied in both the phases discussed here. The site allows C...D nonbonded contacts within the range normally associated with N–D...π interactions, which accounts for its occupation by  $\text{ND}_3$  even when the molecule does not coordinate to a cation. However, although the fully occupied  $\text{ND}_3$  site in the disordered  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$  material (formed by desorption from  $(\text{ND}_3)_8\text{Na}_2\text{C}_{60}$ ) remains in a similar position in  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$ , the subtle change in the number of  $\text{ND}_3$  molecules sorbed by  $\text{Na}_2\text{C}_{60}$  in the two samples, which may reflect the different preparative methods used



**Figure 8.** (a) Sodium site in  $(\text{NH}_3)_{5.3}\text{Na}(\text{C}_{60})$ . The sodium cations are shown as black spheres, the N atoms of the  $\text{NH}_3$  molecules are light gray, and the carbon atoms are medium gray spheres. (b) Comparison of the location of the refined, fully occupied sodium site (black sphere) in  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  with the disordered  $\frac{1}{4}$  occupied sodium sites (small gray spheres) from the parent  $(\text{ND}_3)_6\text{Na}_2\text{C}_{60}$ . (c) Indicates the location of the occupied Na site shown in (b) in the unit cell relative to the  $\text{C}_{60}$  anion (large sphere) on the  $1\text{a}$  position. The disordered sodium sites from the  $Im\bar{3}$  parent structure surrounding each occupied site are represented as light gray spheres. Bold lines connect the Na cation in (b) to its  $\text{C}_{60}$  neighbors.

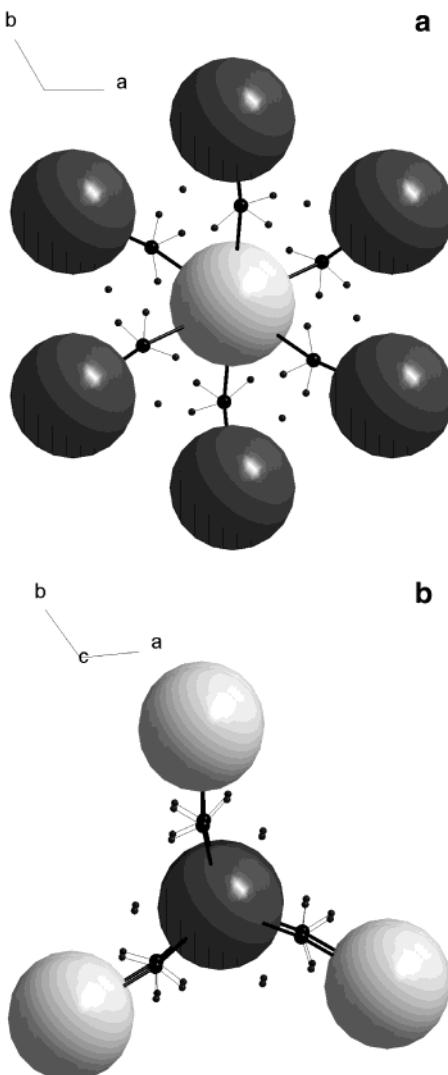
(significant hysteresis is observed in room-temperature isotherm measurements of  $\text{NH}_3$  uptake by  $\text{Na}_2\text{C}_{60}$ ) leads to dramatic differences in the cation site occupancy and fulleride anion orientational ordering. As the disordered  $\text{Na}_2(\text{ND}_3)_6\text{C}_{60}$  is closely related to previously well-studied examples, the discussion below focuses on the new positionally ordered phase.

The major change in the structure of the ordered phase is that instead of the sodium cations occupying a range of sites with 16% probability, one set of these sites is selected for complete occupancy. The ordering of the sodium positions results in retention of only one of the cubic 3-fold axes, and hence drives the lowering of symmetry from cubic to trigonal. The resulting cell can be derived initially by descending in symmetry to rhombohedral, which generates a cell containing three crystallographically equivalent fulleride anions. How-

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**Figure 9.** (a) Arrangement of sodium (black spheres) and nitrogen (medium-sized spheres) around the  $C_{60}$  anions at the  $1a$  sites (large light gray spheres), viewed along the 3-fold axis. (b) Arrangement of sodium and nitrogen around the  $C_{60}$  anions at the  $2d$  sites (large dark gray spheres), viewed along the 3-fold axis. The shorter sodium-centroid distances are directed out of the page.

ever, the cation site ordering requires the loss of the rhombohedral centring translations and results in the development of two distinct fulleride anion sites within the now trigonal cell. The cation site ordering and anion site inequivalence are coupled together in an interesting and potentially significant way. The sodium– $C_{60}$  distances indicate that this differentiation is associated with the development of slightly different charges on the fulleride anions, which also differ in their rotational setting angles about the trigonal [001] direction. The broadening of the  $^{13}\text{C}$  MAS NMR resonance compared with the  $\text{Na}_2\text{C}_{60}$  parent phase is consistent with a more

heterogeneous charge distribution over the carbon atoms in the  $(\text{NH}_3)_{5.3}\text{Na}_2\text{C}_{60}$  phase, although the broadening can also be partially ascribed to  $^{13}\text{C}$ – $^1\text{H}$  dipolar coupling.

The indications of similarities between these materials and the highly correlated transition metal oxide systems are strong. The observation of an electronically interacting array of fulleride anions with crystallographically distinguishable components is significant given the importance of real-space valency ordering in determining the properties of the transition metal oxides, which are archetypal examples of strongly correlated electronic systems. Charge-ordering in oxides is associated with electron localization and insulating behavior.<sup>21</sup> The fullerides are also narrow band systems, as the on-site interelectron repulsion  $U$  is comparable to the  $t_{1u}$  bandwidth  $W$ , with estimates of  $U/W$  placing the  $\text{A}_3\text{C}_{60}$  systems in the localized regime of conventional singly degenerate Mott–Hubbard models. However, in cubic  $\text{A}_3\text{C}_{60}$  systems, the  $t_{1u}$  orbital degeneracy stabilizes the metallic states by opening up alternative hopping channels.<sup>22</sup> Loss of the  $t_{1u}$  degeneracy by symmetry lowering stabilizes magnetic insulating ground states in orthorhombic fullerides<sup>2,3</sup> with the same anion density as metallic cubic phases.<sup>23</sup> In the present case, the interanion separations are similar to those found in magnetic insulating noncubic fullerides.  $(\text{ND}_3)_{5.3}\text{Na}_2\text{C}_{60}$ , however, differs from the previous expanded low-symmetry fullerides in that its structure results from ordering of the charge-compensating sodium cations. This ordering occurs to generate two quite distinct fulleride anion environments, suggesting that it is coupled with localization of the  $t_{1u}$  electrons on the distinct anion sites in an inhomogeneous manner, producing the first example of a charge-ordered fulleride. This is supported by the breadth of the  $^{13}\text{C}$  MAS NMR resonance and the distinct interfulleride contacts at the two sites.

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**Supporting Information Available:** ORTEP and other graphical representations of the structures of the studied compounds (PDF). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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